

**EJERCICIO (08:36)**

Diagonalizar el operador Hamiltoniano de Dirac:

$$\hat{H} = \int d^3x \hat{\psi}^\dagger (-i\gamma^0 \gamma^a \partial_a + \gamma^0 m) \hat{\psi} = \int \frac{d^3p}{(2\pi)^3} \sum_{r=1}^2 E_p \left( \mathbf{a}_{r(p)}^\dagger \mathbf{a}_{r(p)} - \mathbf{d}_{r(p)} \mathbf{d}_{r(p)}^\dagger \right)$$

Donde

$$\hat{\psi} = \sum_{r=1}^2 \int \frac{d^3p}{(2\pi)^3} \frac{1}{\sqrt{2E_p}} \left( \hat{\mathbf{a}}_{r(p)} \mathbf{u}_{r(p)} e^{-ipx} + \hat{\mathbf{d}}_{r(p)}^\dagger \mathbf{v}_{r(p)} e^{ipx} \right)$$

$$\hat{\psi}^\dagger = \sum_{r=1}^2 \int \frac{d^3p}{(2\pi)^3} \frac{1}{\sqrt{2E_p}} \left( \hat{\mathbf{a}}_{r(p)}^\dagger \mathbf{u}_{r(p)}^\dagger e^{ipx} + \hat{\mathbf{d}}_{r(p)} \mathbf{v}_{r(p)}^\dagger e^{-ipx} \right)$$

En el capítulo 8 del curso de Javier habíamos visto que, para los espinores del electrón:

$$\begin{pmatrix} p_0 - m & -\vec{\sigma} \cdot \vec{p} \\ \vec{\sigma} \cdot \vec{p} & -p_0 - m \end{pmatrix} \mathbf{u}_{r(p)} = 0$$

$$\begin{pmatrix} p_0 - m & -\vec{\sigma} \cdot \vec{p} \\ \vec{\sigma} \cdot \vec{p} & -p_0 - m \end{pmatrix} = \begin{pmatrix} p_0 & 0 \\ 0 & -p_0 \end{pmatrix} - \begin{pmatrix} m & 0 \\ 0 & m \end{pmatrix} + \begin{pmatrix} 0 & -\vec{\sigma} \cdot \vec{p} \\ \vec{\sigma} \cdot \vec{p} & 0 \end{pmatrix} = p_0 \gamma^0 - m + \left( \begin{matrix} 0 & -\sigma^a p_a \\ \sigma^a p_a & 0 \end{matrix} \right)$$

$$\begin{pmatrix} p_0 - m & -\vec{\sigma} \cdot \vec{p} \\ \vec{\sigma} \cdot \vec{p} & -p_0 - m \end{pmatrix} = p_0 \gamma^0 - m + \gamma^a p_a$$

$$(p_0 \gamma^0 - m + \gamma^a p_a) \mathbf{u}_{r(p)} = 0$$

$$(-m + \gamma^a p_a) \mathbf{u}_{r(p)} + p_0 \gamma^0 \mathbf{u}_{r(p)} = 0$$

$$(-\gamma^a p_a + m) \mathbf{u}_{r(p)} = p_0 \gamma^0 \mathbf{u}_{r(p)}$$

Multiplicando por  $\gamma^0$  a ambos lados de la ecuación:

$$[1] \quad (-\gamma^0 \gamma^a p_a + \gamma^0 m) \mathbf{u}_{r(p)} = p_0 \mathbf{u}_{r(p)} = E_p \mathbf{u}_{r(p)}$$

Para los espinores del positrón también es válida:

$$\begin{pmatrix} p_0 - m & -\vec{\sigma} \cdot \vec{p} \\ \vec{\sigma} \cdot \vec{p} & -p_0 - m \end{pmatrix} \mathbf{v}_{r(p)} = 0$$

Que resulta en:

$$(-\gamma^0 \gamma^a p_a + \gamma^0 m) \mathbf{v}_{r(p)} = p_0 \mathbf{v}_{r(p)}$$

Pero en este caso, como la energía sería negativa, para lograr un valor “físico” (cap. 8, curso QED de Javier, 10:27) se cambia el signo de la energía:

$$[2] \quad (-\gamma^0 \gamma^a p_a + \gamma^0 m) \mathbf{v}_{r(p)} = -E_p \mathbf{v}_{r(p)}$$

También sabemos por el principio de correspondencia (cap. 1 del curso, minuto 9:06) que  $p_a = i\partial_a$  que aplicando en [1] y [2] resulta:

$$[3] \quad (-i\gamma^0\gamma^a\partial_a + \gamma^0m)u_{r(p)} = E_p u_{r(p)}$$

$$[4] \quad (-i\gamma^0\gamma^a\partial_a + \gamma^0m)v_{r(p)} = -E_p v_{r(p)}$$

Con esto desarrollamos la segunda parte de la integral de  $\hat{H} = \int d^3x \hat{\psi}^\dagger (-i\gamma^0\gamma^a\partial_a + \gamma^0m)\hat{\psi}$

$$(-i\gamma^0\gamma^a\partial_a + \gamma^0m)\hat{\psi} = (-i\gamma^0\gamma^a\partial_a + \gamma^0m) \sum_{r=1}^2 \int \frac{d^3p}{(2\pi)^3} \frac{1}{\sqrt{2E_p}} (\hat{a}_{r(p)} u_{r(p)} e^{-ipx} + \hat{d}_{r(p)}^\dagger v_{r(p)} e^{ipx})$$

Podemos conmutar los operadores con los espinores porque, como dice Javier en los comentarios de video del capítulo "si los objetos vivieran en el mismo espacio, efectivamente habría que cambiar el orden. Sin embargo, los operadores de creación/destrucción actúan sobre el espacio de Hilbert de los estados de base  $|n\rangle$  y los espinores "viven" en el espacio interno (de 4 dimensiones). Es decir, "viven" en espacios diferentes, por tanto conmutan".

$$\begin{aligned} & (-i\gamma^0\gamma^a\partial_a + \gamma^0m)\hat{\psi} \\ &= \sum_{r=1}^2 \int \frac{d^3p}{(2\pi)^3} \frac{1}{\sqrt{2E_p}} \left( (-i\gamma^0\gamma^a\partial_a + \gamma^0m)u_{r(p)} \hat{a}_{r(p)} e^{-ipx} \right. \\ & \quad \left. + (-i\gamma^0\gamma^a\partial_a + \gamma^0m)v_{r(p)} \hat{d}_{r(p)}^\dagger e^{ipx} \right) \end{aligned}$$

Aplicando [3] y [4]:

$$(-i\gamma^0\gamma^a\partial_a + \gamma^0m)\hat{\psi} = \sum_{r=1}^2 \int \frac{d^3p}{(2\pi)^3} \frac{1}{\sqrt{2E_p}} \left( E_p u_{r(p)} \hat{a}_{r(p)} e^{-ipx} - E_p v_{r(p)} \hat{d}_{r(p)}^\dagger e^{ipx} \right)$$

Y reemplazando en el Hamiltoniano:

$$\hat{H} = \int d^3x \hat{\psi}^\dagger (-i\gamma^0\gamma^a\partial_a + \gamma^0m)\hat{\psi}$$

$$\hat{H} = \int d^3x \sum_{s=1}^2 \int \frac{d^3k}{(2\pi)^3} \frac{1}{\sqrt{2E_k}} (\hat{a}_{s(k)}^\dagger u_{s(k)}^\dagger e^{ikx} + \hat{d}_{s(k)} v_{s(k)}^\dagger e^{-ikx}) \sum_{r=1}^2 \int \frac{d^3p}{(2\pi)^3} \frac{E_p}{\sqrt{2E_p}} (u_{r(p)} \hat{a}_{r(p)} e^{-ipx} - v_{r(p)} \hat{d}_{r(p)}^\dagger e^{ipx})$$

$$\hat{H} = \int \frac{d^3k}{(2\pi)^3} \frac{1}{\sqrt{2E_k}} \int \frac{d^3p}{(2\pi)^3} \frac{E_p}{\sqrt{2E_p}} \sum_{s=1}^2 \sum_{r=1}^2 \int d^3x (\hat{a}_{s(k)}^\dagger u_{s(k)}^\dagger e^{ikx} + \hat{d}_{s(k)} v_{s(k)}^\dagger e^{-ikx}) (u_{r(p)} \hat{a}_{r(p)} e^{-ipx} - v_{r(p)} \hat{d}_{r(p)}^\dagger e^{ipx})$$

Trabajamos sobre la integral del espacio:

$$\int d^3x (\hat{a}_{s(k)}^\dagger u_{s(k)}^\dagger e^{ikx} + \hat{d}_{s(k)} v_{s(k)}^\dagger e^{-ikx}) (u_{r(p)} \hat{a}_{r(p)} e^{-ipx} - v_{r(p)} \hat{d}_{r(p)}^\dagger e^{ipx}) =$$

$$\begin{aligned}
 &= \int d^3x \left( \hat{a}_s^\dagger(k) u_s^\dagger(k) e^{ikx} u_{r(p)} \hat{a}_{r(p)} e^{-ipx} + \hat{d}_{s(k)} v_s^\dagger(k) e^{-ikx} u_{r(p)} \hat{a}_{r(p)} e^{-ipx} \right. \\
 &\quad \left. - \hat{a}_s^\dagger(k) u_s^\dagger(k) e^{ikx} v_{r(p)} \hat{d}_{r(p)}^\dagger e^{ipx} - \hat{d}_{s(k)} v_s^\dagger(k) e^{-ikx} v_{r(p)} \hat{d}_{r(p)}^\dagger e^{ipx} \right) = \\
 &= \int d^3x \left( \hat{a}_s^\dagger(k) u_s^\dagger(k) u_{r(p)} \hat{a}_{r(p)} e^{-i(p-k)x} + \hat{d}_{s(k)} v_s^\dagger(k) u_{r(p)} \hat{a}_{r(p)} e^{-i(p+k)x} \right. \\
 &\quad \left. - \hat{a}_s^\dagger(k) u_s^\dagger(k) v_{r(p)} \hat{d}_{r(p)}^\dagger e^{i(p+k)x} - \hat{d}_{s(k)} v_s^\dagger(k) v_{r(p)} \hat{d}_{r(p)}^\dagger e^{-i(k-p)x} \right) = \\
 &= \int d^3x \left( \hat{a}_s^\dagger(k) u_s^\dagger(k) u_{r(p)} \hat{a}_{r(p)} e^{-i(p-k)x} \right) + \int d^3x \left( \hat{d}_{s(k)} v_s^\dagger(k) u_{r(p)} \hat{a}_{r(p)} e^{-i(p+k)x} \right) \\
 &\quad - \int d^3x \left( \hat{a}_s^\dagger(k) u_s^\dagger(k) v_{r(p)} \hat{d}_{r(p)}^\dagger e^{i(p+k)x} \right) - \int d^3x \left( \hat{d}_{s(k)} v_s^\dagger(k) v_{r(p)} \hat{d}_{r(p)}^\dagger e^{-i(k-p)x} \right) =
 \end{aligned}$$

Recordando que  $i(p+k)x = i((E_p + E_k)t - (\vec{p} + \vec{k})\vec{x})$

$$\begin{aligned}
 &= \int d^3x \left( \hat{a}_s^\dagger(k) u_s^\dagger(k) u_{r(p)} \hat{a}_{r(p)} e^{-i(E_p - E_k)t + i(\vec{p} - \vec{k})\vec{x}} \right. \\
 &\quad + \int d^3x \left( \hat{d}_{s(k)} v_s^\dagger(k) u_{r(p)} \hat{a}_{r(p)} e^{-i(E_p + E_k)t + i(\vec{p} + \vec{k})\vec{x}} \right) \\
 &\quad - \int d^3x \left( \hat{a}_s^\dagger(k) u_s^\dagger(k) v_{r(p)} \hat{d}_{r(p)}^\dagger e^{i(E_p + E_k)t - i(\vec{p} + \vec{k})\vec{x}} \right) \\
 &\quad \left. - \int d^3x \left( \hat{d}_{s(k)} v_s^\dagger(k) v_{r(p)} \hat{d}_{r(p)}^\dagger e^{-i(E_k - E_p)t + i(\vec{k} - \vec{p})\vec{x}} \right) = \right. \\
 &= e^{-i(E_p - E_k)t} \hat{a}_s^\dagger(k) u_s^\dagger(k) u_{r(p)} \hat{a}_{r(p)} \int d^3x \left( e^{i(\vec{p} - \vec{k})\vec{x}} \right) \\
 &\quad + e^{-i(E_p + E_k)t} \hat{d}_{s(k)} v_s^\dagger(k) u_{r(p)} \hat{a}_{r(p)} \int d^3x \left( e^{i(\vec{p} + \vec{k})\vec{x}} \right) \\
 &\quad - e^{i(E_p + E_k)t} \hat{a}_s^\dagger(k) u_s^\dagger(k) v_{r(p)} \hat{d}_{r(p)}^\dagger \int d^3x \left( e^{-i(\vec{p} + \vec{k})\vec{x}} \right) \\
 &\quad - e^{-i(E_k - E_p)t} \hat{d}_{s(k)} v_s^\dagger(k) v_{r(p)} \hat{d}_{r(p)}^\dagger \int d^3x \left( e^{i(\vec{k} - \vec{p})\vec{x}} \right) = \\
 &= e^{-i(E_p - E_k)t} \hat{a}_s^\dagger(k) u_s^\dagger(k) u_{r(p)} \hat{a}_{r(p)} (2\pi)^3 \delta^3(p-k) + e^{-i(E_p + E_k)t} \hat{d}_{s(k)} v_s^\dagger(k) u_{r(p)} \hat{a}_{r(p)} (2\pi)^3 \delta^3(p+k) \\
 &\quad - e^{i(E_p + E_k)t} \hat{a}_s^\dagger(k) u_s^\dagger(k) v_{r(p)} \hat{d}_{r(p)}^\dagger (2\pi)^3 \delta^3(p+k) \\
 &\quad - e^{-i(E_k - E_p)t} \hat{d}_{s(k)} v_s^\dagger(k) v_{r(p)} \hat{d}_{r(p)}^\dagger (2\pi)^3 \delta^3(k-p) =
 \end{aligned}$$

Reemplazamos en el Hamiltoniano:

$$\begin{aligned}
 \hat{H} &= \int \frac{d^3k}{(2\pi)^3} \frac{1}{\sqrt{2E_k}} \int \frac{d^3p}{(2\pi)^3} \frac{E_p}{\sqrt{2E_p}} \sum_{s=1}^2 \sum_{r=1}^2 \left( e^{-i(E_p - E_k)t} \hat{a}_s^\dagger(k) u_s^\dagger(k) u_{r(p)} \hat{a}_{r(p)} (2\pi)^3 \delta^3(p-k) \right. \\
 &\quad + e^{-i(E_p + E_k)t} \hat{d}_{s(k)} v_s^\dagger(k) u_{r(p)} \hat{a}_{r(p)} (2\pi)^3 \delta^3(p+k) \\
 &\quad - e^{i(E_p + E_k)t} \hat{a}_s^\dagger(k) u_s^\dagger(k) v_{r(p)} \hat{d}_{r(p)}^\dagger (2\pi)^3 \delta^3(p+k) \\
 &\quad \left. - e^{-i(E_k - E_p)t} \hat{d}_{s(k)} v_s^\dagger(k) v_{r(p)} \hat{d}_{r(p)}^\dagger (2\pi)^3 \delta^3(k-p) \right)
 \end{aligned}$$

Analizamos cada término, cambiando de orden las integrales:

$$\begin{aligned} & \int \frac{d^3p}{(2\pi)^3} \frac{E_p}{\sqrt{2E_p}} \int \frac{d^3k}{(2\pi)^3} \frac{1}{\sqrt{2E_k}} \sum_{s=1}^2 \sum_{r=1}^2 e^{-i(E_p-E_k)t} \hat{a}_s^\dagger(k) u_s^\dagger(k) u_r(p) \hat{a}_r(p) (2\pi)^3 \delta^3(p-k) \\ &= \int \frac{d^3p}{(2\pi)^3} \frac{E_p}{\sqrt{2E_p}} \frac{1}{(2\pi)^3} \frac{1}{\sqrt{2E_p}} \sum_{s=1}^2 \sum_{r=1}^2 e^{-i(E_p-E_p)t} \hat{a}_s^\dagger(p) u_s^\dagger(p) u_r(p) \hat{a}_r(p) (2\pi)^3 \\ &= \int \frac{d^3p}{(2\pi)^3} \frac{E_p}{\sqrt{2E_p}} \frac{1}{\sqrt{2E_p}} \sum_{s=1}^2 \sum_{r=1}^2 \hat{a}_s^\dagger(p) u_s^\dagger(p) u_r(p) \hat{a}_r(p) \end{aligned}$$

$$\begin{aligned} & \int \frac{d^3p}{(2\pi)^3} \frac{E_p}{\sqrt{2E_p}} \int \frac{d^3k}{(2\pi)^3} \frac{1}{\sqrt{2E_k}} \sum_{s=1}^2 \sum_{r=1}^2 e^{-i(E_p+E_k)t} \hat{d}_s(k) v_s^\dagger(k) u_r(p) \hat{a}_r(p) (2\pi)^3 \delta^3(p+k) \\ &= \int \frac{d^3p}{(2\pi)^3} \frac{E_p}{\sqrt{2E_p}} \frac{1}{(2\pi)^3} \frac{1}{\sqrt{2E_{-p}}} \sum_{s=1}^2 \sum_{r=1}^2 e^{-i(E_p+E_{-p})t} \hat{d}_s(-p) v_s^\dagger(-p) u_r(p) \hat{a}_r(p) (2\pi)^3 \\ &= \int \frac{d^3p}{(2\pi)^3} \frac{E_p}{\sqrt{2E_p}} \frac{1}{\sqrt{2E_{-p}}} \sum_{s=1}^2 \sum_{r=1}^2 e^{-i(E_p+E_{-p})t} \hat{d}_s(-p) v_s^\dagger(-p) u_r(p) \hat{a}_r(p) \end{aligned}$$

$$\begin{aligned} & \int \frac{d^3p}{(2\pi)^3} \frac{E_p}{\sqrt{2E_p}} \int \frac{d^3k}{(2\pi)^3} \frac{1}{\sqrt{2E_k}} \sum_{s=1}^2 \sum_{r=1}^2 e^{i(E_p+E_k)t} \hat{a}_s^\dagger(k) u_s^\dagger(k) v_r(p) \hat{d}_r^\dagger(p) (2\pi)^3 \delta^3(p+k) \\ &= \int \frac{d^3p}{(2\pi)^3} \frac{E_p}{\sqrt{2E_p}} \frac{1}{(2\pi)^3} \frac{1}{\sqrt{2E_{-p}}} \sum_{s=1}^2 \sum_{r=1}^2 e^{i(E_p+E_{-p})t} \hat{a}_s^\dagger(-p) u_s^\dagger(-p) v_r(p) \hat{d}_r^\dagger(p) (2\pi)^3 \\ &= \int \frac{d^3p}{(2\pi)^3} \frac{E_p}{\sqrt{2E_p}} \frac{1}{\sqrt{2E_{-p}}} \sum_{s=1}^2 \sum_{r=1}^2 e^{i(E_p+E_{-p})t} \hat{a}_s^\dagger(-p) u_s^\dagger(-p) v_r(p) \hat{d}_r^\dagger(p) \end{aligned}$$

$$\begin{aligned} & \int \frac{d^3p}{(2\pi)^3} \frac{E_p}{\sqrt{2E_p}} \int \frac{d^3k}{(2\pi)^3} \frac{1}{\sqrt{2E_k}} \sum_{s=1}^2 \sum_{r=1}^2 e^{-i(E_k-E_p)t} \hat{d}_s(k) v_s^\dagger(k) v_r(p) \hat{d}_r^\dagger(p) (2\pi)^3 \delta^3(k-p) \\ &= \int \frac{d^3p}{(2\pi)^3} \frac{E_p}{\sqrt{2E_p}} \frac{1}{(2\pi)^3} \frac{1}{\sqrt{2E_p}} \sum_{s=1}^2 \sum_{r=1}^2 e^{-i(E_p-E_p)t} \hat{d}_s(p) v_s^\dagger(p) v_r(p) \hat{d}_r^\dagger(p) (2\pi)^3 \\ &= \int \frac{d^3p}{(2\pi)^3} \frac{E_p}{\sqrt{2E_p}} \frac{1}{\sqrt{2E_p}} \sum_{s=1}^2 \sum_{r=1}^2 \hat{d}_s(p) v_s^\dagger(p) v_r(p) \hat{d}_r^\dagger(p) \end{aligned}$$

Volviendo al Hamiltoniano:

$$\begin{aligned} \hat{H} &= \int \frac{d^3p}{(2\pi)^3} \frac{E_p}{\sqrt{2E_p}} \frac{1}{\sqrt{2E_p}} \sum_{s=1}^2 \sum_{r=1}^2 \hat{a}_s^\dagger(p) u_s^\dagger(p) u_r(p) \hat{a}_r(p) \\ &+ \int \frac{d^3p}{(2\pi)^3} \frac{E_p}{\sqrt{2E_p}} \frac{1}{\sqrt{2E_{-p}}} \sum_{s=1}^2 \sum_{r=1}^2 e^{-i(E_p+E_{-p})t} \hat{d}_s(-p) v_s^\dagger(-p) u_r(p) \hat{a}_r(p) \\ &- \int \frac{d^3p}{(2\pi)^3} \frac{E_p}{\sqrt{2E_p}} \frac{1}{\sqrt{2E_{-p}}} \sum_{s=1}^2 \sum_{r=1}^2 e^{i(E_p+E_{-p})t} \hat{a}_s^\dagger(-p) u_s^\dagger(-p) v_r(p) \hat{d}_r^\dagger(p) \\ &- \int \frac{d^3p}{(2\pi)^3} \frac{E_p}{\sqrt{2E_p}} \frac{1}{\sqrt{2E_p}} \sum_{s=1}^2 \sum_{r=1}^2 \hat{d}_s(p) v_s^\dagger(p) v_r(p) \hat{d}_r^\dagger(p) \end{aligned}$$

Del formulario de Crul, cap. 43, se obtiene:

(<https://crul.github.io/CursoTeoriaCuanticaDeCamposJavierGarcia/#capitulo-43>)

$$u_s^\dagger(p) u_{r(p)} = 2E_p \delta_{rs} \quad (\text{ver fórmula 43.5 y explicación cap. 9 minuto 10:46})$$

$$v_s^\dagger(-p) u_{r(p)} = 0 \quad (\text{ver fórmula 43.5 y ejercicio 43.6.c})$$

$$u_s^\dagger(-p) v_{r(p)} = 0 \quad (\text{ver fórmula 43.5 y ejercicio 43.6.c})$$

$$v_s^\dagger(p) v_{r(p)} = 2E_p \delta_{rs} \quad (\text{ver fórmula 43.5 y explicación cap. 9 minuto 10:46})$$

$$\begin{aligned} \hat{H} = & \int \frac{d^3p}{(2\pi)^3} \frac{E_p}{\sqrt{2E_p}} \frac{1}{\sqrt{2E_p}} \sum_{s=1}^2 \sum_{r=1}^2 \hat{a}_{s(p)}^\dagger 2E_p \delta_{rs} \hat{a}_{r(p)} + 0 - 0 \\ & - \int \frac{d^3p}{(2\pi)^3} \frac{E_p}{\sqrt{2E_p}} \frac{1}{\sqrt{2E_p}} \sum_{s=1}^2 \sum_{r=1}^2 \hat{d}_{s(p)} 2E_p \delta_{rs} \hat{d}_{r(p)}^\dagger \end{aligned}$$

$$\hat{H} = \int \frac{d^3p}{(2\pi)^3} \frac{E_p}{\sqrt{2E_p}} \frac{1}{\sqrt{2E_p}} 2E_p \sum_{s=1}^2 \sum_{r=1}^2 \delta_{rs} \left( \hat{a}_{s(p)}^\dagger \hat{a}_{r(p)} - \hat{d}_{s(p)} \hat{d}_{r(p)}^\dagger \right)$$

$$\boxed{\hat{H} = \int \frac{d^3p}{(2\pi)^3} E_p \sum_{r=1}^2 \left( \hat{a}_{r(p)}^\dagger \hat{a}_{r(p)} - \hat{d}_{r(p)} \hat{d}_{r(p)}^\dagger \right)}$$